

# Rényi entropy and conformal defects

based on 1511.06713 with M. Meineri, R. Myers, M. Smolkin.

Lorenzo Bianchi

Universität Hamburg

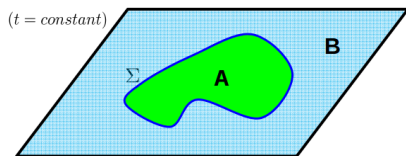


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# Rényi entropy

## Motivation

- It is an **interesting** observable and it is **hard** to compute it in QFT.
- In the limit  $n \rightarrow 1$  gives the **entanglement entropy**.
- It appears in many different contexts: quantum information, condensed matter,...



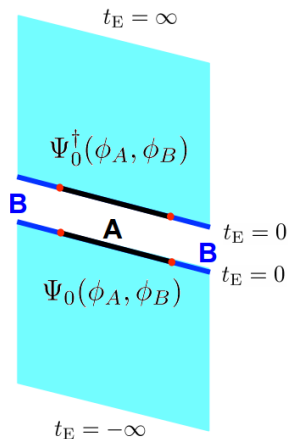
$$\rho_A = \text{Tr}_B(\rho)$$

$$S_n = \frac{1}{1-n} \log \text{Tr}(\rho_A^n)$$

## Entanglement entropy

$$\lim_{n \rightarrow 1} S_n = S_{EE} = -\text{Tr}(\rho_A \log \rho_A)$$

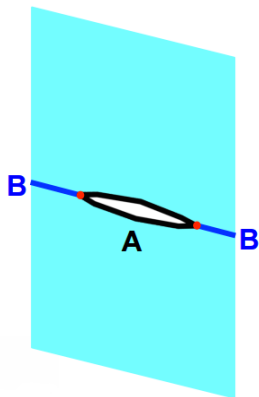
- Path-integral representation of ground state wave-function  $\Psi_0(\phi_A, \phi_B)$



## Rényi entropy

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- Tracing over  $\phi_B$  to obtain the reduced density matrix  $\rho_A$

$$\rho_A(\phi_A^+, \phi_A^-) = \text{Tr}_{\phi_B} \Psi(\phi_A^-, \phi_B) \Psi^\dagger(\phi_A^+, \phi_B)$$



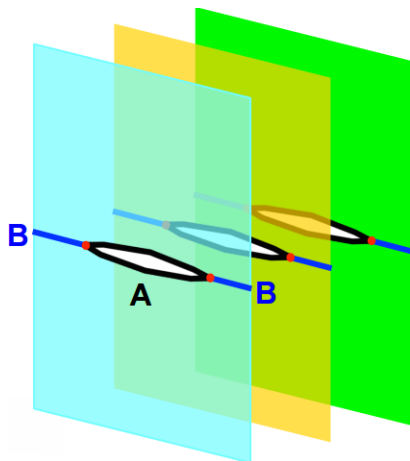
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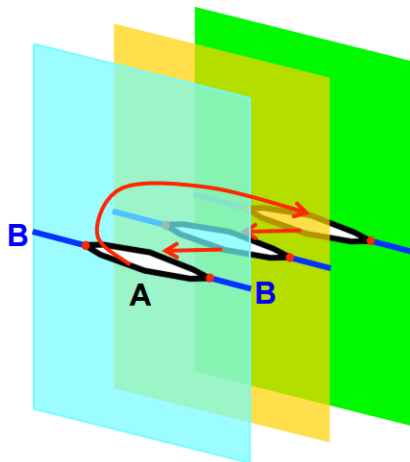
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$$\text{Tr} \rho_a^n = \frac{Z_n}{Z_1^n}$$

- $Z_n$  is the partition function evaluated over a  $n$ -sheeted surface  $\rightarrow$  **HARD**



## Replica trick

Move the problem to target space

$$Z_n = \int [\mathcal{D}\phi]_{\mathcal{R}} \exp \left[ \int_{\mathcal{R}} d^2x \mathcal{L}[\phi](x, t) \right]$$

↓

$$Z_n = \int_{\mathcal{C}_{u,v}} [\mathcal{D}\phi_1 \dots \mathcal{D}\phi_n] \exp \left[ S^{(n)}[\phi_1, \dots, \phi_n] \right]$$

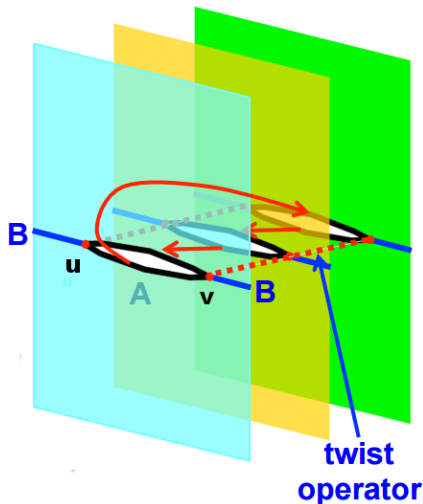
- Restricted path integral

$$\phi_i(x, 0^+) = \phi_{i+1}(x, 0^-) \quad x \in [u, v]$$

- $n$  copies of the same theory

$$S^{(n)}[\phi_1, \dots, \phi_n] = S[\phi_1] + \dots + S[\phi_n]$$

$$S[\phi_i] = \int_{\mathcal{C}} d^2x \mathcal{L}[\phi_i](x)$$



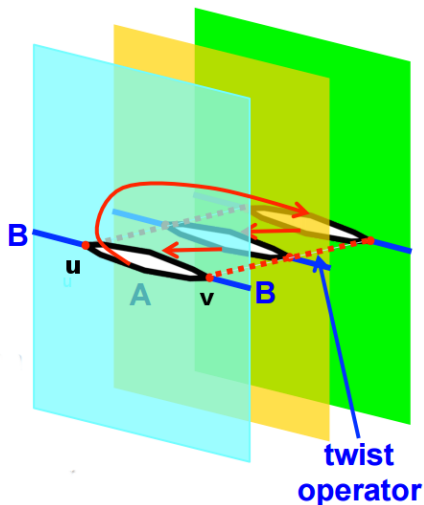
# Rényi entropy

## Twist operators

Location of the cut meaningless: **local effect**



Insertion of local operators in  $u$  and  $v$





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## Higher dimensions

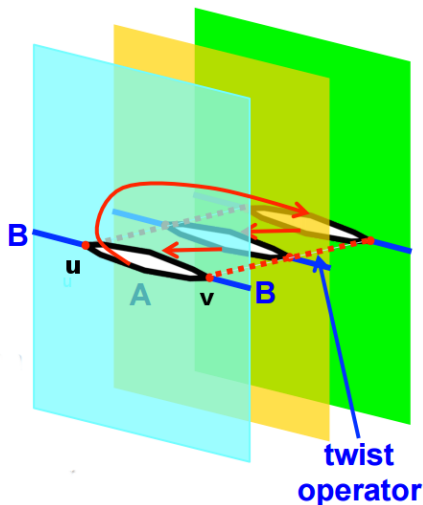
Codimension-two **extended operator** defined on the boundary of  $A$ .

$$\langle \mathcal{T}_n \rangle \equiv \frac{Z_n}{Z_1^n} = e^{(1-n)S_n}.$$

Preserved symmetry

$$SO(d-1, 1) \times U(1)_n$$

**Conformal defect**

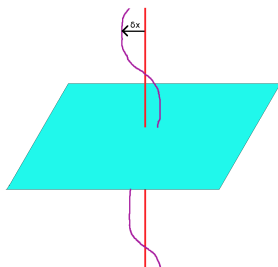


- The defect breaks translation invariance

$$\sum_{m=1}^n \partial_{\mu} T_{(m)}^{\mu a}(x^{\nu}) = \delta_{\Sigma}(x) D^a(x^i),$$

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- $D^a(x^i)$  is the **displacement operator**
- It implements small modifications of the defect

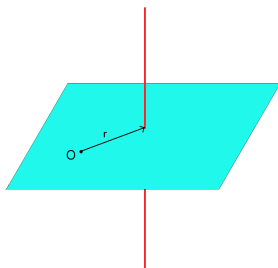
$$\delta \langle X \rangle_n = - \int d^{d-2}x \delta x^a(x^i) \langle D^a(x^i) X \rangle_n$$

- Its two-point function is fixed by symmetry

$$\langle D^a(x^i) D^b(0) \rangle_n = C_D(n) \frac{\delta^{ab}}{|x^i|^{2(d-1)}}.$$

- The defect breaks translation invariance

$$\sum_{m=1}^n \partial_\mu T_{(m)}^{\mu a}(x^\nu) = \delta_\Sigma(x) D^a(x^i),$$



- Local operators acquire a non-vanishing **one-point function**.
- The kinematics is fixed by conformal invariance

$$\langle O \rangle_n \equiv \frac{\langle \tau_n O \rangle}{\langle \tau_n \rangle} = \frac{C_O}{r^\Delta}$$

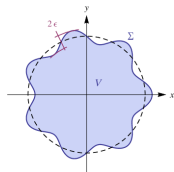
- For the stress tensor

$$\langle T_{ij} \rangle_n = -\frac{h_n}{2\pi n} \frac{\delta_{ij}}{r^d} \quad \langle T_{ab} \rangle_n = \frac{h_n}{2\pi n} \frac{(d-1)\delta_{ab} - d n_a n_b}{r^d}$$

## Three different conjectures

- Entropy across a deformed sphere [Mezei, 2014]

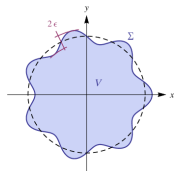
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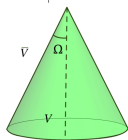
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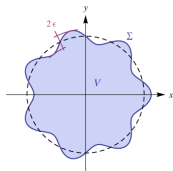
$$S_n^{\text{univ}} \stackrel{\Omega \rightarrow \pi/2}{\sim} \begin{cases} a(d) \frac{h_n}{n(n-1)} (\Omega - \frac{\pi}{2})^2 \log(\ell/\delta) & d \text{ odd} \\ b(d) \frac{h_n}{n(n-1)} (\Omega - \frac{\pi}{2})^2 \log^2(\ell/\delta) & d \text{ even} \end{cases}$$



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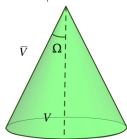
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- Entropy in 4d for arbitrary smooth entangling surface [Lee, McGough, Safdi, 2014]

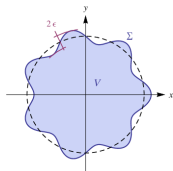
$$S_n = \left( -\frac{f_a(n)}{2\pi} \int_{\partial A} R_\Sigma - \frac{f_b(n)}{2\pi} \int_{\partial A} \tilde{K}_{ij}^a \tilde{K}_{ij}^a + \frac{f_c(n)}{2\pi} \int_{\partial A} \gamma^{ij} \gamma^{kl} C_{ikjl} \right) \log(\mu\ell) + \lambda_n,$$

$$f_b(n) = f_c(n)$$

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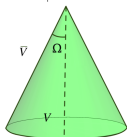
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$$f_b(n) = f_c(n)$$

Equivalent to: [LB, Meineri, Myers, Smolkin, 2014]

$$C_D(n) = d \Gamma\left(\frac{d+1}{2}\right) \left(\frac{2}{\sqrt{\pi}}\right)^{d-1} h_n \quad \partial_n C_D|_{n=1} = d \Gamma\left(\frac{d+1}{2}\right) \left(\frac{2}{\sqrt{\pi}}\right)^{d-1} \partial_n h_n|_{n=1} = \frac{2\pi^2}{d+1} C_T$$



- We developed a **field theoretic framework** for calculating the dependence of Rényi entropies on the **shape of the entangling surface** in a conformal field theory.
- We **unified different conjectures**, which were proven for the entanglement entropy [Faulkner, Leigh, Parrikar, 2015], but recently seemed to fail holographically for any  $n$  [Dong, 2016].
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THANK YOU